

The lattice of signed measures

Notation If ν is a positive measure and f is ν -integrable, denote by ν_f the signed measure satisfying

$$\int y \, d\nu_f = \int y f \, d\nu$$

for all ν -integrable y .

We need a simple measure-theoretic fact:

Prop. $\nu_f \geq 0$ if and only if $f \geq 0$ ν -a.e.

PR The reverse direction (\Leftarrow) is obvious, so let us prove the forward direction (\Rightarrow), or rather its contrapositive.

Suppose $f < 0$ on a non-null set. Then

$$0 < \nu \{ f < 0 \} = \nu \left(\bigcup_{n \in \mathbb{N}} \{ f < -\frac{1}{n} \} \right)$$

$$\leq \sum_{n \in \mathbb{N}} \nu \{ f < -\frac{1}{n} \}$$

Given $\epsilon > 0$:

so there exists $\epsilon := \gamma_n > 0$ such that
 $\mu\{f < -\epsilon\} > 0$.

Conclude that

$$\begin{aligned}\mu_f\{f < -\epsilon\} &= \int_{f < -\epsilon} f \, d\mu \\ &\leq -\epsilon \int_{f < -\epsilon} d\mu \\ &< 0.\end{aligned}$$

□

Cor For $f, g \in L^1(\mu) = L^1(\mu; \mathbb{R})$,

$\mu_f \leq \mu_g$ iff $f \leq g$ μ -a.e.

PF. Apply the proposition to $\mu_{f-g} = \mu_f - \mu_g$.

□

Let $\mathcal{M}(\mu)$ denote the space of signed measures absolutely continuous w.r.t. μ . By the Radon-Nikodym theorem, the

map $L^1(\mu) \rightarrow \mathcal{M}(\mu)$, $f \mapsto \mu_f$

is a vector space isomorphism. If we recall that the total variation of a signed measure $\nu \ll \mu$ is defined (independently of μ , it turns out) by $|\nu| = \mu_{|f|}$ for $f \in L^1(\mu)$ satisfying $\nu = \mu_f$, then we can define the norm of ν to be

$$\|\nu\| := |\nu|(X) = \int |f| \, d\mu = \|f\|_{L^1(\mu)}.$$

Thus the map $f \mapsto \mu_f$ exhibits an isomorphism of $L^1(\mu)$ with $\mathcal{M}(\mu)$ as normed vector spaces and even as Banach spaces:

$$L^1(\mu) \cong \mathcal{M}(\mu) \text{ in Ban.}$$

All this is standard textbook material in real analysis.

What is equally obvious, but less commonly remarked on,

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is that (as the corollary shows), the map $f \mapsto \mu_f$ makes $L^1(\omega)$ and $M(\omega)$ isomorphic as posets:

$$L^1(\omega) \cong M(\omega) \text{ in Pos.}$$

\uparrow pointwise order μ -a.e. \uparrow setwise order

Cor. $M(\omega)$ is a lattice, with meets and joins given by

$$\mu_f \wedge \mu_g = \mu_{f \wedge g}$$

$$\mu_f \vee \mu_g = \mu_{f \vee g}$$

Pr. $L^1(\omega)$ has meets and joins, given by

$$f \wedge g = \min(f, g)$$

$$f \vee g = \max(f, g) \quad \text{pointwise.}$$

□

In fact, the lattice structure extends to the space of all signed measures. Let $M(X)$ denote the space of signed measures on a measurable space X , so that for any positive measure ν on X , $M(\omega) \in M(X)$.

Thm $M(X)$ is a lattice, with meets and joins given by

$$\nu \wedge \rho = \nu \frac{d\nu}{d\nu} \wedge \frac{d\rho}{d\nu}, \quad \nu \vee \rho = \nu \frac{d\nu}{d\nu} \vee \frac{d\rho}{d\nu},$$

where ν is any ^{positive} measure dominating ν and ρ .

Pr. First, note that such a measure ν always exists; simply take $\nu = |\nu| + |\rho|$. We must also show that the RHSs are well-defined.

Suppose we have two positive measures ν_1, ν_2 with $\nu, \rho \ll \nu_1, \nu_2$.

Put $\nu = \nu_1 + \nu_2$, and calculate

$$\left(\frac{d\nu}{d\nu_1} \wedge \frac{d\rho}{d\nu_1} \right) d\nu_1 \stackrel{\text{chain rule}}{=} \left(\frac{d\nu}{d\nu_1} \wedge \frac{d\rho}{d\nu_1} \right) \frac{d\nu_1}{d\nu} d\nu$$

$$\stackrel{\text{positivity}}{=} \left(\frac{d\nu}{d\nu_1} \frac{d\nu_1}{d\nu} \wedge \frac{d\rho}{d\nu_1} \frac{d\nu_1}{d\nu} \right) d\nu$$

$$\stackrel{\text{chain rule}}{=} \left(\frac{d\nu}{d\nu} \wedge \frac{d\rho}{d\nu} \right) d\nu.$$

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This proves that the quantity defining meets is well-defined, and similarly for joins.

It remains to prove that the quantity actually is a meet.

Let $\sigma \in \mathcal{M}(X)$ be a lower bound for ν, ρ : $\sigma \leq \nu$ and $\sigma \leq \rho$.

Put $\omega = |\nu| + |\rho| + |\sigma|$, s. that $\nu, \rho, \sigma \in \mathcal{M}(\omega)$. By what we have proved earlier, $\nu \wedge \rho$ is (well-)defined above is a meet (greatest lower bound) in $\mathcal{M}(\omega)$, hence $\sigma \leq \nu \wedge \rho$.

Since $\sigma \in \mathcal{M}(X)$ is arbitrary, $\nu \wedge \rho$ is also a meet in $\mathcal{M}(X)$.

The argument for joins is dual. \square

We will use the lattice structures of signed measures, $\mathcal{M}(X)$, and of positive measures, $\mathcal{M}^+(X)$, to define a composition of subcoupling measures.

Remark One might naively try to define the meet of measures ~~with~~ ν, ρ by

$$(\nu \wedge \rho)(A) \stackrel{?}{:=} \nu(A) \wedge \rho(A).$$

However, the resulting "measure" need not even be additive, hence need not be a measure at all.